First Revision exam – Jan 2025 Chengalpattu District XII - Physics Answer Key

$$
V_{\text{galvanometer}} = V_{\text{shunt}}
$$

$$
\Rightarrow I_g R_g = (I - I_g)S
$$

$$
S = \frac{I_g}{(I - I_g)} R_g \text{ or }
$$

$$
I_g = \frac{S}{S + R_g} I
$$

$$
\theta = \frac{1}{G} I_g \Rightarrow \theta \propto I_g \Rightarrow \theta \propto I \text{ So,}
$$

$$
\frac{1}{R_{\text{eff}}} = \frac{1}{R_{\text{g}}} + \frac{1}{S} \Rightarrow R_{\text{eff}} = \frac{R_{\text{g}}S}{R_{\text{g}} + S} = R_{\text{a}}
$$

 $S = \frac{f_x}{(I - I_x)} R_g$ or
 $I_x = \frac{S}{S + R_g} I$

Since, the deficition in the galvanometer

is proportional to the current passing

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galvanometer is a m

$$
\frac{\Delta I}{I}{\times}100\%=\frac{I_{\text{ideal}}-I_{\text{actual}}}{I_{\text{ideal}}}{\times}100\%
$$

32 Capacitor not only stores the charge but also it stores energy. When a battery is connected to the capacitor, electrons of total charge –*Q* are transferred from one plate to the other plate. To transfer the charge, work is done by the

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The electric field at a point C due to +*q* is
\n
$$
\vec{E}_{+} = \frac{1}{4\pi\epsilon_0} \frac{q}{(r-a)^2}
$$
 along BC

Since the electric dipole moment vector \vec{p} is from $-q$ to $+q$ and is directed along BC,

$$
\vec{E}_{+} = \frac{1}{4\pi\epsilon_0} \frac{q}{(r-a)^2} \hat{p}
$$
 (1.13)

$$
\vec{E}_{-} = -\frac{1}{4\pi\epsilon_0} \frac{q}{(r+a)^2} \hat{p}
$$
 (1.14)

$$
\vec{E}_{tot} = \vec{E}_{+} + \vec{E}_{-}
$$
\n
$$
= \frac{1}{4\pi\epsilon_{0}} \frac{q}{(r-a)^{2}} \hat{p} - \frac{1}{4\pi\epsilon_{0}} \frac{q}{(r+a)^{2}} \hat{p}
$$

the above equation is rewritten as
\n
$$
\frac{\overline{E}_{+} = \frac{1}{4\pi\epsilon_{0}} \frac{q}{(r-a)^{2}} \hat{P}
$$
\n(1.13)
\nwhere \hat{P} is the electric dipole moment unit
\nvector from -q to +q.
\nThe electric field at a point C due to -q is
\n
$$
\frac{\overline{E}}{\overline{E}_{-} = -\frac{1}{4\pi\epsilon_{0}} \frac{q}{(r+a)^{2}} \hat{P}
$$
\n(1.14)
\nSince +q is located closer to the point C than
\n-q, \overline{E}_{+} is stronger than
\nlength of the \overline{E}_{+} vector is drawn larger than
\nthat of \overline{E}_{+} vector is drawn larger than
\nthat of \overline{E}_{-} vector.
\nThe total electric field at point C is
\ncalculated using the superposition principle
\nof the electric field.
\n
$$
\overline{E}_{tot} = \frac{q}{4\pi\epsilon_{0}} \left(\frac{1}{(r-a)^{2}} - \frac{1}{(r+a)^{2}} \right) \hat{P}
$$
\n(1.15)
\n
$$
\overline{E}_{tot} = \frac{1}{4\pi\epsilon_{0}} \left(\frac{4\pi a}{(r-a)^{2}} - \frac{1}{(r+a)^{2}} \right) \hat{P}
$$
\nNote that the total is along
\n \vec{E}_{+} , since +q is closer to C than -q. The
\ndirection of \overline{E}_{tot} is shown in Figure 1.16.
\n
$$
\frac{A}{Q} = \frac{1}{\epsilon_{0}} \left(\frac{4\pi a}{(r-a)^{2}} \right) \hat{P}
$$
\n(1.16)
\n
$$
\frac{A}{Q} = \frac{1}{\epsilon_{0}} \left(\frac{4\pi a}{(r-a)^{2}} \right) \hat{P} + \frac{\overline{E}_{\omega}}{C}
$$
\nIf the point C is very far away from the
\ndipole $(r > \epsilon_{0})$. Then under this limit the
\nterm $(r^{2} - a^{2})^{2} = r^{4}$. Substituting this into
\nequation (1.16), we get
\n
$$
\overline{E}_{tot} = \frac{4\pi(\epsilon_{0}}{4\pi(\epsilon_{0})^{2}} \hat{P} (r > > a)
$$

$$
\begin{array}{c|c}\nA & a & B \\
\hline\nq & 0 & q \\
\hline\nr & & r\n\end{array}
$$

$$
\overline{\vec{E}_{tot}} = \frac{1}{4\pi\epsilon_0} \left(\frac{4aq}{r^3}\right) \hat{p} \quad (r > > a)
$$
\n
$$
\text{since } 2aq \,\hat{p} = \bar{p}
$$
\n
$$
\vec{E}_{tot} = \frac{1}{4\pi\epsilon_0} \frac{2\,\vec{p}}{r^3} \quad (r > > a) \tag{1.17}
$$

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Wavefronts from S_1 and S_2 spread out and overlap on the other side of the double slit. When a screen is placed at a distance of about 1 m from the slits, alternate bright and dark fringes which are equally spaced appear on the second Theorem . There are called the control of the second temperature where the second temperature of α the second temperature of α the second of α the second of α the second of α the second

$$
\beta = y_{(n+1)} - y_n = \left[(n+1) \frac{\lambda D}{d} \right] - \left[n \frac{\lambda D}{d} \right]
$$

$$
\beta \text{ for bright, } \beta = \frac{\lambda D}{d} \tag{7.31}
$$

$$
\beta = y_{(n+1)} - y_n = \left(\frac{(2(n+1)-1)\lambda D}{2}\right) - \left(\frac{(2n-1)\lambda D}{2}\right)
$$

$$
\beta \text{ for dark, } \beta = \frac{\lambda D}{d}
$$
 (7.32)

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$$
d\vec{B} = \frac{\mu_0}{4\pi} \frac{Idl \sin \theta}{r^2} i
$$

First equation

It is nothing but the Gauss's law of electricity. It relates the net electric flux to net electric charge enclosed in a surface. Mathematically, it is expressed as

 $\oint_C E \cdot d\lambda = \frac{Q_{\text{train}}}{c}$
 \downarrow

(Semistric field and Q enclosed

where \vec{E} is the electric field and Q enclosed

where \vec{E} is the electric field and Q enclosed.

This equation is true for both discurses

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Third equation

It is Faraday's law of electromagnetic induction. This law relates electric field with the changing magnetic flux which is mathematically written as

$$
\oint_{l} \vec{E} \cdot d\vec{l} = -\frac{d}{dt} \Phi_{B} \quad \text{(Faraday's law)} \tag{5.9}
$$

where \vec{E} is the electric field. This equation implies that the line integral of

$$
\oint_{l} \vec{B}.\vec{dl} = \mu_{0}i_{C} + \mu_{0}\epsilon_{0}\frac{d}{dt}\oint_{s}\vec{E}.\vec{dA}
$$
\n(Ampere-Maxwell law)

\n
$$
(5.10)
$$

equations are known as Maxwell's equations in electrodynamics. This where \vec{B} is the magnetic field. This equation shows that both conduction current and displacement current produce magnetic field. These four equation ensures the existence of electromagnetic waves. The entire communication system in the world depends on electromagnetic waves. In fact our understanding of stars, galaxy, planets etc come by analysing the electromagnetic waves emitted by these astronomical objects. the electric field around any closed path

is equal to the rate of change of magnetic

flux through the closed path bounded by

the surface.

Our modern technological revolution

to the to Franchy have of electromagnetic

36. b) **Microscopic model of current**

$$
\begin{array}{c}\n \begin{array}{c}\n \stackrel{dx}{\longrightarrow} \\
 \hline\n \stackrel{0}{\longrightarrow} \\
 \hline\n \stackrel{0}{\longrightarrow} \\
 \hline\n \stackrel{0}{\longrightarrow} \\
 \hline\n u_d \stackrel{dt}{\longrightarrow}\n \end{array}\n \end{array}
$$

The drift velocity of the electrons = v_A If the electrons move through a distance dx within a small interval of dt , then

$$
\nu_{d} = \frac{dx}{dt}; \quad dx = \nu_{d} dt \qquad (2.7)
$$

$$
= A \, dx \times n \tag{2.8}
$$

$$
= (A \, v_{\scriptscriptstyle A} \, dt \,) \, n
$$

$$
dQ = (e)(Avddt)n
$$

Hence the current $I = \frac{a}{2}$

$$
J = \frac{1}{A}
$$

$$
J = nev
$$

Since A is the area of cross section of
the conduction the chief contexts available in the
two conducts the decisions available in the
values α wanter of decisions per unit volume
 $\alpha = A\alpha x/n$
 $= A\alpha x/n$
Solution give tran

As a result of flux change, emf is induced in both primary and secondary coils. The emf induced in the primary coil or back emf ε _o is given by

$$
\epsilon_{\rm p}=-N_{\rm p}\frac{d\Phi_{\rm B}}{dt}
$$

$$
\upsilon_{p} = -N_{p} \frac{d\Phi_{B}}{dt}
$$
\n(4.24)

$$
\varepsilon_{s} = -N_{s} \frac{d\Phi_{B}}{dt}
$$

$$
v_s = -N_s \frac{d\Phi_B}{dt} \tag{4.25}
$$

$$
\frac{\nu_s}{\nu_p} = \frac{N_s}{N_p} = K \tag{4.26}
$$

$$
\frac{\nu_s}{\nu_p} = \frac{N_s}{N_p} = \frac{i_p}{i_s}
$$

interference of electrons diffracted from various atomic layers of the target material. From the known value of interplanar spacing of Nickel, the wavelength of the electron wave was experimentally calculated as 1.65Å. The wavelength can also be calculated from de Broglie relation for $V =$ 54 V $\frac{N = \frac{N}{N} + \frac{N}{N} + \frac{N}{N}}{N}$

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